初期の物理

(CGC, Glasma, 流体化や熱化) と課題について

along along

福嶋健二 (東京大学)

高度化後のALICE実験での物理の可能性

Q1: What is "Early Thermalization"?

What is "Early Thermalization"

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Color Glass Condensate (CGC) $\tau \lesssim 1/Q_s \sim 0.1 {\rm fm/c}$

Color Glass + Plasma = Glasma $\tau \lesssim \tau_0 \sim \Lambda_{\rm QCD}^{-1}$

(s) Quark-Gluon Plasma $\tau \lesssim au_f \sim 10 {
m fm/c}$

Hadronization (quarks \rightarrow hadrons)

What is "Early Thermalization" , Mersel, Mersel, Mersel, MerselsMersel, Mersel, Mersel, Mersel, Me Can Thermalization in Heavy Ion Collisions be

Described by QCD Diagrams?

Yuri V. Kovchegov^{*}

pQCD expert

Department of Physics, The Ohio State University Columbus. OH 43210

March, 2005

and pioneer of AdS/CFT for HIC

CGC : $1/\tau \rightarrow \text{Bjorken} : 1/\tau ^{4/3}$ **How?** Impossible?

What is "Early Thermalization"

algos algos

Fukushima, RoPP (2016)



Three concepts clearly distinguished since Chesler-Yaffe (2010)







a symptotically AdS_5 spacetime

A 02139, USA*

attle, WA 98195, USA^{\dagger}

One of the most important holographic papers



Valid ity of hydro does not require thermalization nor isotropization!

Not known in other fields

Q2: Physics of isotropization?



Color Glass Condensate (CGC)

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Stasto-Golec-Biernat-Kwiecinski Plot

Geometric Scaling

Scaling variable

$$Q_s^2(x) = Q_0^2 (x/x_0)^{-\lambda}$$

CGC already seen? Extended geometric scaling

Color Glass Condensate (CGC)

Finding CGC =

Finding Scaling Properties (with Qs)

Possible in ep, pA, etc, but not so clear in AA

Color Glass Condensate (CGC)

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Finding CGC =

Finding necessity of special class of resummation (Wilson lines)

Will be discussed later

Glasma

Early-time dynamics of heavy-ion collision described in the CGC approximation



Initial longitudinal pressure is **negative** Isotropic plasma has positive pressure

Glasma

Gelis-Epelbaum, PRL (2013), see also a plenary talk @ QM2013



However, UV contamination unavoidable

Glasma

allow allow

Maybe $P_L / P_T \sim 0.6$ or smaller (free streaming?)

Q: Why not very small ? A: Very viscous ! → aHydro

Please tell me what can be measured to probe it ?

I asked this question to Berndt 4 years ago, and his answer was "balance function."

Any observable to quantify anisotoropization ? Anything sensitive to initial negative P_L ?

Q3: Physics of hydrodynamization?

Hydrodynamization

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Fukushima, RoPP (2016) based on Strickland



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Hydrodynamization

allow allow

$$P_{\rm T} - P_{\rm L} = \frac{2\eta}{\tau}$$

or equivalently

$$\frac{P_{\rm L}}{P_{\rm T}} = \frac{3\tau T - 16(\eta/s)}{3\tau T + 8(\eta/s)}$$

 $P_{\rm L}/P_{\rm T} \approx 0.5$ for initial $\tau_0 \sim 0.5 \, {\rm fm}/c$ and $T_0 \sim 0.4 \, {\rm GeV}$, if $\eta/s \sim 1/(4\pi)$

Hydrodynamization

Definition of "Hydrodynamization" is unclear

Hydrodynamization occurs for viscous hydro far earlier than ideal hydro...

Hydrodynamization is a theoretical concept I do not think there is any observable for this...

For more details please ask Hirano-san, Nonaka-san

Q4 : Physics of thermalization?

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Distribution function of gluons

Local thermal eq. :
$$f(p) \sim \frac{1}{e^{p \cdot u(x)/T(x)} - 1}$$

Time evolution ?
CGC initial condition : $f(p) \sim p^{-\nu}$
perturbative tail

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Fukushima, RoPP (2016) based on Berges



Fixed-points of the Boltzmann equations ~ Power law (Kolmogorov-Zakharov spectrum)

Initial-time numerical simulation is reliably done as long as the spectrum is power-law.

UV divergence if the spectrum is power-law.

No simulation valid over the whole momentum region

Dense gluon

small-p

Classical statistical sim.

Dilute gluon

Boltzmann equations

large-p

KZ Spectrum ~ (Wave) Turbulence ~ Non-thermal Fixed-Point

Overpopulation (2011)

Bose–Einstein Condensation and Thermalization of the Quark Gluon Plasma

Jean-Paul Blaizot⁽¹⁾, François Gelis⁽¹⁾, Jinfeng Liao⁽²⁾, Larry McLerran^(2,3), Raju Venugopalan⁽²⁾

Observable ? Some value of v

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Is "turbulence" observable?

In principle yes : Early-time gluons with power-law spectrum may affect hadron spectrum and correlation

Caveats :

- * Perturbative calculations → Power-law
- * Hard to tell power-law from exponential

CGC in pA (or forward AA)

Flow Observables

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n-th harmonics of *m*-particles (nucl-th/0105040)

$$\kappa_n\{m\} := \prod_{i=1}^m \int \frac{\mathrm{d}^2 \boldsymbol{p}_{\perp i}}{(2\pi)^2} \,\mathrm{e}^{\mathrm{i}n(-1)^{i+1}\phi_i} \,\frac{\mathrm{d}^m N}{\mathrm{d}^2 \boldsymbol{p}_{\perp 1} \cdots \boldsymbol{p}_{\perp m}}$$

In the dipole model



Flow Observables

Fukushima-Hidaka (2017) motivated by Dusling-Mace-Venugopalan (2017)



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Flow Observables

Recent developments along the same lines

Mace-Skokov-Tribedy-Venugopalan (2018)

$$v_{2n}^2\{2\} \propto N_{\rm ch}^0 = v_{2n+1}^2\{2\} \propto N_{\rm ch}$$

CGC-based prediction not depending on a dipole model

Systematics can tell the CGC-type correlation from hydro

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Benic-Fukushima-Garcia-Montero-Venugopalan (2018)



Photon in pp at LHC

Good agreement! More soft photons needed! Many "CGC" calculations assume the *k_T* factorization

Compton Scattering

 $q \xrightarrow{\gamma}$

Annihilation



 $\propto \alpha_e \alpha_s n_q (1 - n_q) n_g \propto \alpha_e \alpha_s n_q n_{\bar{q}} (1 + n_q)$ $(qg \rightarrow q\gamma)$ $(q\bar{q} \rightarrow q\gamma)$

Compton Scattering

 $q - \frac{q}{g}$

$$\propto \alpha_e \alpha_s n_q (1 - n_q) \alpha_s^{-1}$$

~ $\alpha_e n_q (1 - n_q)$

Annihilation





Gelis-Jalilian-Marian (2002)



 $\sim \alpha_e n_q \langle UU^{\dagger}$

$$\rho_p$$
 ρ_A

$$\sim \alpha_e n_q n_{\bar{q}} \langle U U^{\dagger} U U^{\dagger} \rangle$$

Annihilation process suppressed by quark distribution

Multiple Scattering with CGC

CGC Calculation for Photon $\frac{1}{A_{\perp}} \frac{d\sigma^{q \to q\gamma}}{d^2 \boldsymbol{k}_{\perp}} = \frac{2\alpha_e}{(2\pi)^4 \boldsymbol{k}_{\perp}^2} \int_0^1 dz \frac{1 + (1-z)^2}{z} \int d^2 \boldsymbol{l}_{\perp} \frac{\boldsymbol{l}_{\perp}^2 C(\boldsymbol{l}_{\perp})}{(\boldsymbol{l}_{\perp} - \boldsymbol{k}_{\perp}/z)^2}$ $C(\boldsymbol{l}_{\perp}) \equiv \int d^2 \boldsymbol{x}_{\perp} e^{i\boldsymbol{l}_{\perp} \cdot \boldsymbol{x}_{\perp}} e^{-B_2(\boldsymbol{x}_{\perp})} = \int d^2 \boldsymbol{x}_{\perp} e^{i\boldsymbol{l}_{\perp} \cdot \boldsymbol{x}_{\perp}} \left\langle U(0)U^{\dagger}(\boldsymbol{x}_{\perp})\right\rangle_{\rho}$ $B_2(\boldsymbol{x}_\perp - \boldsymbol{y}_\perp) ~\equiv~ Q_s^2 \int d^2 \boldsymbol{z}_\perp [G_0(\boldsymbol{x}_\perp - \boldsymbol{z}_\perp) - G_0(\boldsymbol{y}_\perp - \boldsymbol{z}_\perp)]^2$ 0.3 0.25 **GC** Effect 0.2 seen at soft C(k) 0.15 momentum 0.1 + crossed diagram 0.05 (photon emitted first) ρ_{A} $\rho_{\rm p}$ 0 0 5 10 15 20 25 30 k/Λ_{OCD} Gelis-Jalilian-Marian (2002) 34 Aug. 18, 2018 @ Nagasaki

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CGC Calculation for Photon LO: ~ $\alpha_e n_q \langle UU^{\dagger} \rangle$ **NLO:** $\sim \alpha_e \langle (g\rho_p)^2 \rangle \langle UU^{\dagger}UU^{\dagger} \rangle$ $(g\rho_p)^2 < n_q \le g\rho_p$

NLO is overwhelming but the pA expansion still works

Systematic calculations feasible Not small corrections but dominant at high energies pA photon data (hopefully) coming very soon

Benic-Fukushima-Garcia-Montero-Venugopalan (2016)



Benic-Fukushima-Garcia-Montero-Venugopalan (2016)





This approximation makes sense when a large momentum (or quark mass) is involved in the considered process

Then, the distribution function is introduced, in which a part of resummation is taken into account



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Gluons are surely dominant degrees of freedom

Heavy flavor is hard to see the true CGC effect

To see the true CGC effect, direct photons ~ a few GeV

Is this possible? How far?

In pA the nuclear PDF could be probed directly

Photons ~ a few GeV is interesting also as a probe to detect strong magnetic fields and induced phenomena